

The role of the virtual work in Faraday's law.

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Abstract

In the context of Faraday's induction law, we show that the concept of virtual work provides another point of view to clarify the nature of the induced electric field, as well as the fact that the integral over a closed path of the induced electric field is not the work performed by a unit charge. The usefulness of the concept of virtual conservativity is discussed. Further, we study the relation between the electromotive force and the real work done by an external agent to keep a circuit at constant velocity. From this discussion it is observed that magnetic forces can be non-conservative from the mathematical point of view, but can be treated as conservative for all practical purposes.

KEYWORDS: Faraday's law, virtual work, virtual conservativity.

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1 Introduction

The central nature of the electrostatic field leads to its conservative character, and its field lines start and finish on charges. Faraday's experiments revealed that electric fields with non-null circulation can be created (in absence of net charge) by time-varying magnetic fluxes. These electric fields have closed field lines, so that they are non-conservative. Faraday's law has a crucial role as a theoretical tool [1]-[5] as well as in applications [6]. Therefore, we encounter several points of view to explain Faraday's Law and the nature of the induced electric field [1]-[10]. Nevertheless, it is not usual to distinguish between the non-conservative behavior of induced fields* with respect to the non-conservative behavior of fields that come from charges in motion. In the same way, no distinction is usually done concerning closed integrals of the field defined in a fixed time, with the real work on a closed path done over a unit charge. In this paper, we analyze the role of the virtual work as an alternative tool to understand Faraday's induction law and to find the nature of the induced electric field, taking into account that the integral over a closed path of such a field is not the work performed by a unit charge. In addition, we study the relation between the electromotive force and the real work done by an external agent to keep a circuit at constant velocity. The latter relation leads us to discuss some subtleties concerning the conservative nature of the magnetic forces.

2 Faraday's induction law

In order to initiate the discussion, we present briefly the customary approach on Faraday's law. Consider a charge q in a rectangular loop of wire that moves at constant velocity \mathbf{v} , in a non-uniform stationary magnetic field \mathbf{B} , (e.g. produced by a long conducting wire, as shown in Fig. 1). We shall calculate the line integral of $\mathbf{f} \cdot d\mathbf{r}$ over a closed path in the rectangular loop, where \mathbf{f} is the force due to the magnetic force on the charge q . In the segment closer to the wire, charges tend to circulate in counter-clockwise sense (observed from above), while in the segment farther to the wire, charges tend to circulate in the clockwise sense as observed from above. Since \mathbf{B} is more intense in the closer segment, net circulation goes in counter-clockwise sense. As for the segments that are parallel to \mathbf{v} , the displacement of q and the force on it are perpendicular, so they give no contribution to $\mathbf{f} \cdot d\mathbf{r}$. Taking the contribution from the segments perpendicular to \mathbf{v} , the so-called electromotive force yields.

$$\varepsilon = \frac{1}{q} \oint \mathbf{f} \cdot d\mathbf{r} = vw (\mathbf{B}_1 - \mathbf{B}_2) \quad (1)$$

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*In this document, the concept of induced electric field will be used with its meaning in the context of Faraday's law.

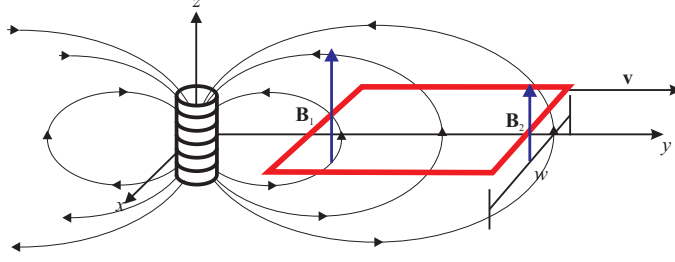


Figure 1: A rectangular loop moving at constant velocity \mathbf{v} , in the field generated by a long conducting wire. \mathbf{B}_1 is the field in the segment of the loop that is closer to the wire, while \mathbf{B}_2 is the field in the farthest segment, w is the width of the loop.

where \mathbf{B}_1 , \mathbf{B}_2 and w are defined in Fig. 1. In textbooks, it is customary to argue that when the charge moves very rapidly, the integral (1) represents the work realized over q [5]. We take a textual fragment from Ref. [5] Sec. 7.3

“If we imagine a charge q to move all around the loop, in a time short enough so that the position of the loop has not changed appreciably, then Eq. (1) gives the work done by the force \mathbf{f} per unit charge”.

Nevertheless, owing to the form of the Lorentz force $q\mathbf{v} \times \mathbf{B}$, the magnetic field cannot do work regardless the velocity of the charge. The contradiction comes from the fact that the line integral in Eq. (1), is carried out at a **fixed time**. Therefore, it does not correspond to a real work, since a real work must consider time-evolution [4, 11]. In our subsequent developments, we shall show the following (a) Equation (1) can also be interpreted as a virtual work done by the magnetic force \mathbf{f} on the charge q . (b) In some cases ε coincides with the real work done by an external agent to keep a circuit at constant velocity[†]. (c) Magnetic forces are not conservative from the mathematical point of view but they can be considered as conservative for all practical purposes.

By now, we come back to the problem of the rectangular loop. Let us consider the flux due to the magnetic field through the rectangular loop. It is independent of the form of the surface limited by the loop because $\nabla \cdot \mathbf{B} = 0$. The change of flux along the time dt gives

$$d\Phi = -vw(\mathbf{B}_1 - \mathbf{B}_2) dt \quad (2)$$

combining Eqs. (1, 2) we find

$$\varepsilon = -\frac{d\Phi}{dt} \quad (3)$$

and this relation is valid for an arbitrary form and velocity of the loop [5, 7]. Moreover, for an observer F' that moves with the loop, the electric and magnetic fields are \mathbf{E}' and \mathbf{B}' respectively. Since for F' the loop is at rest, the electromotive force (emf) comes exclusively from the electric field, and we find

$$\varepsilon' = \oint \mathbf{E}' \cdot d\mathbf{r}' = -\frac{d\Phi'}{dt'} \quad (4)$$

When we extrapolate the previous equations to an arbitrary closed path C which is stationary with respect to an inertial frame F , we obtain Faraday’s induction law

$$\varepsilon = \oint_C \mathbf{E}(\mathbf{r}, t) \cdot d\mathbf{r} = -\frac{d}{dt} \int_S \mathbf{B}(\mathbf{r}, t) \cdot d\mathbf{S} \quad (5)$$

where S is a surface which has C as its boundary. It is essential to emphasize that the closed line integral for the electric field in Eq. (5), is performed for a **fixed value of time**.

3 Virtual work

In order to analyze the role of the virtual work [12, 13] in Faraday’s law, it is important to point out the nature of the electromotive force (emf). Under the conditions of such a law, the emf owes only to the electric field, because the closed path is assumed stationary with respect to the system of reference that measures the flux. A second aspect is that the emf is calculated as an integral over a closed path in which each differential element is calculated at the same value of time t . For example, two contributions in two different segments of the path $\mathbf{E}(x_1, y_1, z_1, t)d\mathbf{r}_1$ and $\mathbf{E}(x_2, y_2, z_2, t)d\mathbf{r}_2$

[†]This item was shown by Refs. [4, 11] but with some mistakes in the procedure. Indeed, the correction of the procedure led us to some subtleties about the conservative nature of the magnetic force.

are evaluated at the same time t . It implies that such an integral is not the work that a point unit charge would do to close the circuit, for if the field is a function of time, a differential of it must be calculated using the value of the field in the space-time point where the particle is located, two small contributions for this real work must be of the form $\mathbf{E}(x_1, y_1, z_1, t_1) d\mathbf{r}_1$ and $\mathbf{E}(x_2, y_2, z_2, t_2) d\mathbf{r}_2$.

On the other hand, to calculate the work done by the magnetic force in Eq. (1), it was assumed that the line integral is realized at a fixed instant of time. We call this a **virtual work**, which is the work done on a charge when time is considered fixed along the path of integration. In the virtual displacement defined here, the only constraint is that the particle must follow the prescribed path C . Notice that it does not correspond to the real work that would be done on a charge, since real works must consider time evolution. In that sense, we can say that magnetic forces cannot do real work, but they can do virtual work. According to this discussion, $d\mathbf{r}$ in Eq. (1) is a virtual displacement.

Further, it is well-known that in the general case, the induced field comes from the relativistic transformation of the fields from the system F to the system F' in which the loop is stationary, where F' is connected with F by a boost. In F' the contribution to the closed line integral comes exclusively from the electric field, and we see that this closed line integral is not zero. Now, since this field comes from a relativistic transformation, it is not generated by charges.

The concept of virtual work permits us to see the same facts under another point of view. It is clear that the electric fields generated by charges in motion are in general non-conservative because of the explicit time-dependence. Nevertheless, these fields are **virtually conservative**. To show this, observe that in a virtual process, owing to the lack of time evolution, the configuration of charges that are the sources of the field are fixed as we travel over the path to calculate the line integral. Therefore, in a virtual process an electric field coming from charges behaves as an electrostatics field, and therefore it is virtually conservative[‡]. By contrast, the induced field involved in Faraday's law is not virtually conservative, since the virtual work (emf) over a closed path is non-zero. It implies that such a field cannot come from charges in motion. In that sense, the closed path integral in Faraday's law is not the correct quantity to evaluate real conservativity, except under certain approximations.

There is still another way to see the difference between fields coming from charges, and induced fields. For fields generated by charges, the condition $\nabla \times \mathbf{E}_{charge} = 0$ is satisfied in the whole space and at any given time, while for induced fields $\nabla \times \mathbf{E}_{induc} \neq 0$ in at least some points of space-time. The condition $\nabla \times \mathbf{E}_{charge} = 0$ is necessary and sufficient for conservativity in electrostatics. Notwithstanding, for charges in motion, in which \mathbf{E} is a function of time, the nullity of the curl only guarantees that $\mathbf{E} = -\nabla\phi(\mathbf{r}, t)$ which in turn implies that a virtual work yields

$$q \int_{r_A}^{r_B} \mathbf{E} \cdot d\mathbf{r} = q [\phi(\mathbf{r}_B, t) - \phi(\mathbf{r}_A, t)]$$

and it is independent of the path since time is fixed in the process. However, in a real trajectory the work can depend on the path because of the time evolution. The conclusion is that $\nabla \times \mathbf{E} = 0$ in the whole space at any time, is equivalent to **virtual conservativity** (i.e. conservativity for virtual works). Of course, in the electrostatic case virtual works coincide with real works so that in this case, virtual conservativity is equivalent to real conservativity.

Despite the induced field \mathbf{E} is not virtually conservative, nor is originated in charges, the force undergone by a charge immersed in this field takes the form $\mathbf{F} = q\mathbf{E}$, which is the expression found by fields generated by charges. So although $\mathbf{E}_{charges}$ and \mathbf{E}_{induc} come from different sources, they accomplish the same local property with which electric field was initially defined. From the previous discussion the total electric field can be decomposed as

$$\mathbf{E} = \mathbf{E}_{charge} + \mathbf{E}_{induc}, \quad (6)$$

where

$$\begin{aligned} \nabla \times \mathbf{E}_{charge} &= 0 \quad ; \quad \nabla \times \mathbf{E}_{induc} = -\frac{\partial \mathbf{B}}{\partial t} \\ \nabla \cdot \mathbf{E}_{charge} &= \frac{\rho}{\varepsilon_0} \quad ; \quad \nabla \cdot \mathbf{E}_{induc} = 0 \end{aligned}$$

So that \mathbf{E}_{charge} is irrotational (i.e. virtually conservative), and its field lines start and finish on charges. By contrast, \mathbf{E}_{induc} is solenoidal (charges are not its sources) and it is not virtually conservative.

4 Electromotive force as a real work

We have related the emf with the concept of virtual work. In this section, we shall see that in some cases the emf can also be related with real works. To show this, let us consider a U-conductor with a conducting rod sliding on it with constant

[‡]Of course there is a (instantaneous) magnetic field generated by the charges in motion. But this emf arises from the electric field only.

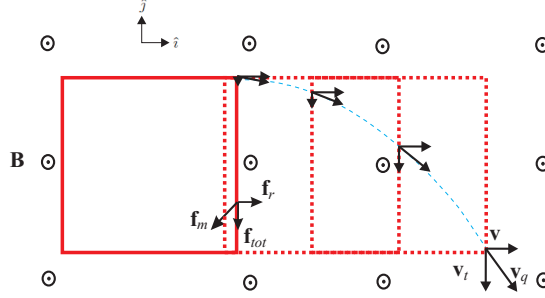


Figure 2: A U-shape conductor with a conducting road sliding on it at constant velocity, by the action of an external agent. The system is immersed in a uniform magnetic field going out of the paper. In our picture, we have assumed $q > 0$, but all equations are valid for an arbitrary value and sign of q .

velocity \mathbf{v} . The system is immersed in a uniform magnetic field $\mathbf{B} = B\hat{\mathbf{k}}$ going out of the paper (see Fig. 2). There are two forces acting on the charge q , the magnetic force \mathbf{f}_m due to \mathbf{B} , and \mathbf{f}_r which is a “drag” force that is provided for an external agent (to keep the velocity of the rod constant), and transmitted to the charge by the structure of the wire[§]. Hence, the total force \mathbf{f}_{tot} on the charge q , yields

$$\mathbf{f}_{tot} = \mathbf{f}_r + \mathbf{f}_m \quad (7)$$

if we denote \mathbf{v}_t the tangential velocity of the charge, the real velocity of the charge is

$$\mathbf{v}_q = \mathbf{v} + \mathbf{v}_t = v\hat{i} + v_t\hat{j}$$

in a similar way a real infinitesimal displacement of the charge $d\mathbf{L}_q$ is given by

$$d\mathbf{L}_q = d\mathbf{L} + d\mathbf{L}_t = \hat{i} dL + \hat{j} dL_t$$

The x -component of the total force on q is zero because \mathbf{v} is constant. Hence

$$\begin{aligned} \mathbf{f}_r \cdot \hat{i} &= -\mathbf{f}_m \cdot \hat{i} = -q [(\mathbf{v} + \mathbf{v}_t) \times \mathbf{B}] \cdot \hat{i} \\ \mathbf{f}_r \cdot \hat{i} &= -qv_t B = f_r \end{aligned} \quad (8)$$

Now, the total force along the y -direction is given by the y -component of the magnetic force

$$f_{tot} = \mathbf{f}_{tot} \cdot \hat{j} = \mathbf{f}_m \cdot \hat{j} = -qvB \quad (9)$$

hence the net force on q is constant and in the y -direction. Therefore, the trajectory of the charge is a parabola (and not a straight line as assumed in Ref. [11]). Calculating the work made by \mathbf{f}_r , we obtain

$$dW_r = \mathbf{f}_r \cdot d\mathbf{L}_q = f_r \hat{i} \cdot (\hat{i} dL + \hat{j} dL_t) = f_r dL = -qBv_t(t) v dt$$

Integrating for a charge starting at the top of the rod and ending at the bottom of it, we have

$$W_r = -qBv \int_{t_0}^{t_f} v_t(t) dt = -qBv \int_{\mathbf{r}_0}^{\mathbf{r}_f} dL_t$$

if $q > 0$, the charge goes downward from top to bottom, hence

$$\int_{\mathbf{r}_0}^{\mathbf{r}_f} dL_t = \int_w^0 dL_t = -w$$

if $q < 0$ this integral is w . Therefore

$$W_r = |q| Bvw$$

Now, since the magnetic force does not work, W_r must coincide with the work done by the total force \mathbf{f}_{tot} as can be checked easily

[§]Microscopically, this drag force is due to a Hall effect inside the wire, see Ref.[11].

$$W_{f_{tot}} = \int_{\mathbf{r}_0}^{\mathbf{r}_f} \mathbf{f}_{tot} \cdot d\mathbf{L}_q = -qvB \int_{\mathbf{r}_0}^{\mathbf{r}_f} dL_t = |q| Bvw$$

further, we can observe that these real works coincide with the virtual work done by the magnetic force, i.e. with the emf

$$W_m^{virt} = \int \mathbf{f}_m \cdot d\mathbf{L}_t = \int_{\mathbf{r}_0}^{\mathbf{r}_f} (\mathbf{f}_m \cdot \hat{j}) dL_t = -qvB \int_{\mathbf{r}_0}^{\mathbf{r}_f} dL_t = |q| Bvw$$

this result coincides with the one found in Ref. [11]. It is understandable since the total force is conservative, so despite Ref. [11] is using a wrong trajectory, the length covered along the y -axis is the same. However, the fact that we have a uniform acceleration instead of a uniform velocity brings an apparent paradox as we discuss below.

Since the motion along the y -axis is accelerated, the speed v_t depends on time. In turn it implies that \mathbf{f}_r depends explicitly on time according with Eq. (8). Further Eq. (9) shows that the total force \mathbf{f}_{tot} is constant. Consequently \mathbf{f}_m must also depend explicitly on time to cancel the time-dependence of \mathbf{f}_r in Eq. (7). But forces that depends explicitly on time are non-conservative, so \mathbf{f}_m is a force that depends explicitly on time but never does real work!.

This paradox has to do with the initial definition of a conservative vector field. A vector field \mathbf{F} is conservative if for **any trajectory**, the line-integral along such a trajectory only depends on the extreme points. In the case of magnetic forces, we only take the path traced by **the real trajectory of the particle**, if we calculated the line integral along other trajectories (with fixed extremes) we would obtain in general a result dependent on the particular trajectory[¶]. Furthermore, if we assume an arbitrary magnetic field and that velocities depend on the position, the explicit calculation of the rotational of the magnetic force gives in general a non-zero value. This is related with the fact that there is no potential function ϕ that depends on the position only, and that generates the magnetic force through the relation $\mathbf{F} = -\nabla\phi$.

To solve the puzzle we simply observe that in Physics we are not concern about arbitrary line-integrals, we only mind line-integrals over real trajectories of the particles to calculate works. When applied to these real trajectories, magnetic forces produce null line-integrals, line-integrals of such forces on other trajectories are just a mathematical curiosity. We could say that magnetic forces are not conservative from the mathematical definition, but they can be treated as conservative for practical purposes. Hence, it is natural to have a magnetic force that depends explicitly on time (as it is our case), but never does a work.

5 Conclusions

The introduction of the concept of virtual work to analyze Faraday's induction law permits to clarify the nature of the induced electric field, which does not come from charges, and the fact that the integral over a closed path of the induced electric field is not the real work that the field would do on a unit charge along the closed path. It also allows to understand why in certain systems of reference, the closed line integral of the magnetic force is not zero, despite magnetic fields cannot do work.

We studied the problem of a U-conductor with a conducting rod sliding on it with constant velocity \mathbf{v} , in which the system is immersed in a uniform magnetic field. We observe that the emf in this system coincides with the real work done by the external agent to keep constant the velocity of the rod. In this example, the magnetic force depends explicitly on time and thus is not conservative as a vector field in the mathematical sense. We observe however, that magnetic forces can be treated as conservative for practical purposes in Physics.

On the other hand, the fact that $\nabla \times \mathbf{E} = 0$ for the whole space at any time, is equivalent to virtual conservativity of \mathbf{E} , that is conservativity in a process that does not involve time-evolution. In general, it differs from real conservativity in which time evolution is essential. Electric fields coming from charges are virtually conservative (and sometimes conservative), while induced electric fields in the Faraday sense are not virtually conservative. Therefore, the concept of virtual conservativity is useful to distinguish between induced electric fields in the Faraday sense, and electric fields generated by charges.

It is worth pointing out that the concept of virtual conservativity can be extended to any vector field in Physics, and therefore could be useful in Mechanics and other scenarios of Physics.

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[¶]Two identical paths define different trajectories if they travel with different time dependence. For vector fields depending explicitly on time, identical paths covered in different time intervals lead to different values of the line integral.

References

- [1] W. Taussig Scott, “*The Physics of Electricity and Magnetism*”. 2nd Ed. New York. John Wiley & Sons, Inc. (1966).
- [2] Gaylord P. Harnwell, “*Principles of Electricity and Electromagnetism*”. 2nd Ed. New York. McGraw-Hill Book Company Inc. (1949).
- [3] J. D. Jackson “*Classical Electrodynamics*” 3rd Ed. New York. John Wiley & Sons. (1998).
- [4] David J. Griffiths “*Introduction to Electrodynamics*”. 3rd Ed. New Jersey, Prentice Hall. (1999).
- [5] Edward M. Purcell, “*Electricity and Magnetism*” Berkeley Physics Course Vol. 2, 2nd Ed. New York. McGraw-Hill International Editions. (1985).
- [6] W. M. Saslow, “*Electromechanical implications of Faraday’s law: A problem collection*” Am. J. Phys. **55** (1987) 986.
- [7] A. López-Ramos, J. R. Menéndez and C. Piqué, “*Conditions for the validity of Faraday’s law of induction and their experimental confirmation*” Eur. J. Phys. **29** (2008) 1069.
- [8] Dragan V. Redžic, “*Various paths to Faraday’s law*” Eur. J. Phys. **29** (2008) 257.
- [9] I. Galili, D. Kaplan and Y. Lehavi, “*Teaching Faraday’s law of electromagnetic induction in an introductory physics course*” Am. J. Phys. **74** (2006) 337.
- [10] Dragan V. Redžic, “*Faraday’s law via the magnetic vector potential*” Eur. J. Phys. **28** (2007) N7.
- [11] E. P. Mosca, “*Magnetic Forces Doing Work?*” Am. J. Phys. **42** (1974) 295.
- [12] S. N. Srivastava, “*Method of Virtual Work for Some Problems in Physics*” Am. J. Phys. **28** (1960) 299.
- [13] W. P. Healy, “*A line-integral work function for the Lorentz force*” Eur. J. Phys. **18** (1997) 350.